An overview of synthetic spectrum generation methods. A look at the Kurucz models

> Urtzi Jauregi Fakulteta za Matematiko in Fiziko FMF Astro debata, 21.04.2009

- We now have databases with hundreds of *thousands* of spectra (RAVE, SDSS), which will soon become hundreds of *millions* (GAIA).
- Manual analysis is impossible.
- Therefore, we need highly efficient techniques to automatically extract physical parameters from observed spectra.
- •The most obvious: compare observations with a database ("grid") of synthetic spectra.

The process of creating spectra



A synthetic spectrum Teff=5500, logg = 4.5, [M/Mo] = 0



Current trends: 1D vs 3D

- 1D models use a precalculated solution of the ERT with tabulated opacities to compute line shape.
- Fast and reliable
- Static
- Don't account for some of the interesting physics: convection (→ errors in Teff) and time-dependent phenomena.
- Need for "fudge" parameters (e.g. for turbulence)

Current trends: 1D vs 3D

- 3D models solve the full hydrodynamic problem (equations of conservation of mass & energy) with tabulated opacities.
- They are better.
- Time-dependent solutions.
- Physics "emerge" naturally, without need for "fudge" parameters.
- Very slow, unsuitable for large-scale work (15 min/line!).
- Highly nonlinear problem.

SST Observation and Numerical Simulation





Large spatial variation of emergent line profiles

Line formation in inhomogeneous stellar atmospheres

Line formation in inhomogeneous stellar atmospheres



The space-time averaged 3D spectrum is NOT the same as the spectrum of the space-time averaged 3D atmosphere ! Line formation *non-linear* ⇒ 3D abundance corrections

1D models: PP vs Spherical

At the core of every 1D model is the Equation of Radiative Transfer:

$$\cos\theta \frac{\partial I_{\nu}}{\partial \tau_{\nu}} - \frac{\sin\theta}{\tau_{\nu}} \frac{\partial I_{\nu}}{\partial \theta} = I_{\nu} - S_{\nu}$$

Spherical approximation

$$\cos\theta \frac{dI_{\nu}}{d\tau_{\nu}} = I_{\nu} - S_{\nu}$$

Plane-parallel approximation (thin photosphere, homogeneous layers)

 $d \tau_v = \kappa_v \rho dr$ is the optical depth of the medium.

PHOENIX (Hautschild)

- Developed by P. Hautschild (H. Sternwarte);
- plane-parallel or spherical, heavily NLTE.
- static or (up to relativistic!) expanding media;
- opacity for ca. 592 million lines (42 atomic, 550 molecular);
- <u>Code not available</u>; must run on his computers.

http://www.hs.uni-hamburg.de/EN/For/ThA/phoenix/index.html



- Developed by B. Gustafsson et al (University of Uppsala);
- plane-parallel or spherical, LTE.
- only pre-calculated models and low-res fluxes! You have to calculate spectra with your own software;
- roughly comparable with Kurucz models;
- <u>Code not available.</u>

The Kurucz "family"

- Developed by R. Kurucz (Harvard) in the early 70s; precursor of several modern methods.
- Plane-parallel, (mostly) LTE models. Many simplifications to optimize code for speed.
- Metallicities scaled to solar (ATLAS 9) or individually changeable (ATLAS 12).
- Large line database (162 million lines!).
- All code and data publicly available. Active community with several support mailing lists.
- Scarce documentation, confusing I/O, non-standard code . . . http://kurucz.harvard.edu/

Integrating the ERT gives us several magnitudes as a function of depth in the photosphere.

Flux:

$$F_{\nu}(\tau_{\nu}) = 2\pi \left[\int_{\tau_{\nu}}^{\infty} S_{\nu} E_{2}(t_{\nu} - \tau_{\nu}) dt_{\nu} - \int_{0}^{\tau_{\nu}} S_{\nu} E_{2}(\tau_{\nu} - t_{\nu}) dt_{\nu}\right]$$

where

$$E_n(x) \equiv \int_1^\infty \frac{e^{-xs}}{s^n} ds$$

is the *exponential integral* function.

Average intensity:

$$J_{\nu}(\tau_{\nu}) = 1/2 \int_{\tau_{\nu}}^{\infty} S_{\nu} E_{1}(t_{\nu} - \tau_{\nu}) dt_{\nu} - 1/2 \int_{0}^{\tau_{\nu}} S_{\nu} E_{1}(\tau_{\nu} - t_{\nu}) dt_{\nu}$$

Radiation pressure:

$$P_{\nu}(\tau_{\nu}) = 1/2 \int_{\tau_{\nu}}^{\infty} S_{\nu} E_{3}(t_{\nu} - \tau_{\nu}) dt_{\nu} - 1/2 \int_{0}^{\tau_{\nu}} S_{\nu} E_{3}(\tau_{\nu} - t_{\nu}) dt_{\nu}$$

The source function is a kernel that contains all the information about the radiation field.

In the presence of both absorption and scattering,

$$S(\tau) = \frac{\kappa^{a}}{\kappa^{s} + \kappa^{a}} B_{T} + \frac{\kappa^{s}}{\kappa^{s} + \kappa^{a}} J(\tau)$$

S depends on J,

but

J depends on S.

Solution: iterative methods. But, how to begin?

Assuming the absorption coefficient κ_{ν} does not depend on ν , the source function can be written as

$$S(\tau) = \frac{3}{4\pi} (\tau + \frac{2}{3}) F(0)$$

and

$$T = \left[\frac{3}{4}(\tau + \frac{2}{3})\right]^{1/4} T_{eff}$$

We need to compute observables such as flux $F_n(t)$. Flux depends on source function $S_n(t)$. So:

- 1. We compute an initial *S* by means of a Gray model;
- 2. we use the result from the ERT to compute average intensity *J*, flux, etc.;
- 3. we use J to re-compute S;
- 4. we repeat the procedure until getting convergence.

In real atmospheres, spectral lines get broadened by different physical phenomena. Kurucz uses:

• Quantum (radiative) broadening: due to the Uncertainty Principle. Affects all lines.

• Stark broadening: Caused by protons and free electrons. Affects most lines specially in hot stars.

Van der Waals broadening. Caused by neutral H.
 Affects most lines in cold stars

Broadening profiles are convolved with each other, the result being a Voigt function H(a,v), where

$$a = \frac{\Gamma_R + \Gamma_S + \Gamma_W}{4 \pi \nu_D}$$

Is the width of the profile.

For each depth, Kurucz assumes that a line profile is given by

$$k_{\nu} = \left[\frac{\pi e^2}{mc} \frac{N}{Z} \frac{1}{\rho} gf \frac{1}{\sqrt{\pi} \Delta v_D}\right] \left[e^{-E/kT}\right] \left[H(a, \nu)\right] \left[1 - e^{-h\nu/kT}\right]$$

All factors after the first are <1, so if the first is <0.001, the line is discarded. If not:

- The line is accepted and the procedure repeated;
- the continuum opacity is calculated and added;
- intensity at every wavelength is calculated, with the ERT.
- We jump to the next depth.

The spectrum is the convolved with a rotation profile and an instrumental one, to account for rotational broadening and limited instrumental resolution. The result is the final spectrum.

Kurucz spectrum vs observed



- Synthetic spectra are a rapidly-evolving tool for an extremely demanding task;
- The future is in full HD 3D ab-initio models; however, 1D will be mostly used for the next few years;
- There is a wide choice of 1D model/spectrum suites, but the ones whose code is public spur the most development.
- We are bound to see a lot happening in this field yet.

_0×

```
COMMON /BAL/BAL1(kw,9), AAL1(kw), SAL1(kw), XNFPAL(kw,2), BAL2(kw,1)
 COMMON /BB/BB1(kw,7),XNFPB(kw,1)
 COMMON /BC/BC1(kw,14),AC1(kw),SC1(kw),XNFPC(kw,2),BC2(kw,6)
 COMMON /BCA/BCA1(kw,8),BCA2(kw,5),XNFPCA(kw,2)
 COMMON /BFE/BFE1(kw,15),AFE1(kw),SFE1(kw),XNFPFE(kw,1)
 COMMON /BHE/BHE1(kw,29),AHE1(kw),SHE1(kw),BHE2(kw,6),AHE2(kw),
             SHE2(kw),AHEMIN(kw),SIGHE(kw),XNFPHE(kw,3),XNFHE(kw,2)
 COMMON /BHYD/BHYD(kw,8),AHYD(kw),SHYD(kw),AH2P(kw),BMIN(kw),
              AHMIN(kw), SHMIN(kw), SIGH(kw), SIGH2(kw), AHLINE(kw),
              SHLINE(kw),XNFPH(kw,2),XNFH(kw)
COMMON /BK/BK1(kw,8),XNFPK(kw,1)
 COMMON /BMG/BMG1(kw,11),AMG1(kw),SMG1(kw),XNFPMG(kw,2),BMG2(kw,6)
 COMMON /BNA/BNA1(kw,8),XNFPNA(kw,1)
COMMON /B0/B01(kw,13),XNFPC(kw,),FOP(kw,4)
COMMON /BSI/BSI1(kw,11),ASI1(kw),SSI (kw,4),XNFPSI(kw,2),BSI2(kw,10)
 COMMON /CHOH/XNFPCH(kw),XNFPOH(kw)
 COMMON /CONT/ABTOTC(kw),ALPHAC(kw),TAUNUC(kw),SNUC(kw),HNUC(kw),
              JNUC(kw), JMINSC(kw), RESIDC(kw)
REAL*8 JNUC, JMINSC
COMMON /CONV/DLTDLP(kw),HEATCP(kw),DLRDLT(kw),VELSND(kw),
              GRDADB(kw), HSCALE(kw), FLXCNV(kw), VCONV(kw), MIXLTH,
2
              IFCONV
REAL*8 MIXLTH
 COMMON /ELEM/ABUND(99),ATMASS(99),ELEM(99),XABUND(99),WTMOLE
 COMMON /FLUX/FLUX,FLXERR(kw),FLXDRV(kw),FLXRAD(kw)
 COMMON /FREQ/FREQ,FREQLG,EHVKT(kw),STIM(kw),BNU(kw),WAVENO
 COMMON /FRESET/FRESET(500),RCOSET(500),NULO,NUHI,NUMNU,IFWAVE,
```