

Peculiar velocities and Type Ia supernovae

Vid Iršič

Advisor: dr. Anže Slosar

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Outline

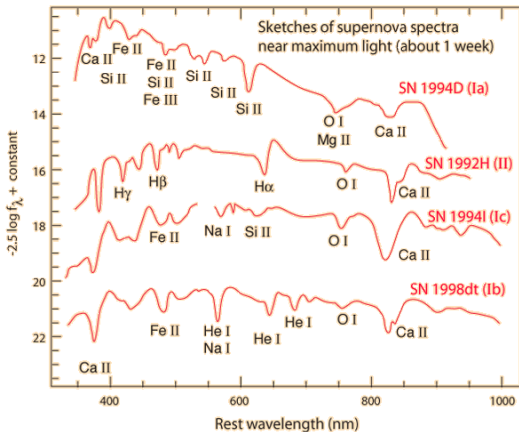
- 1 Overture
- 2 Type Ia supernovae
 - Standard Candles
 - Data corrections
- 3 Distances in cosmology
 - Friedmann equations
 - Luminosity distance
- 4 Cosmological parameters from SNe Ia
 - Peculiar velocities
- 5 Conclusion

Overture

Different types of supernovae (SNe):

- Type II (H lines)

Overture



Sketches of spectra from Carroll & Ostlie, data attributed to Thomas Matheson of National Optical Astronomy Observatory.

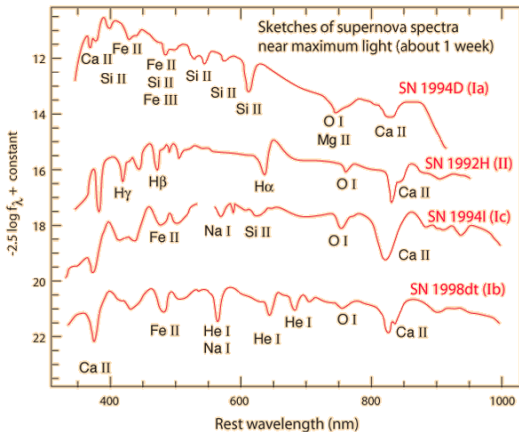
<http://hyperphysics.phy-astr.gsu.edu/hbase/hframe.html> (Dec 18., 2009)

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Different types of supernovae (SNe):

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- Type Ib (no H lines, He I lines)

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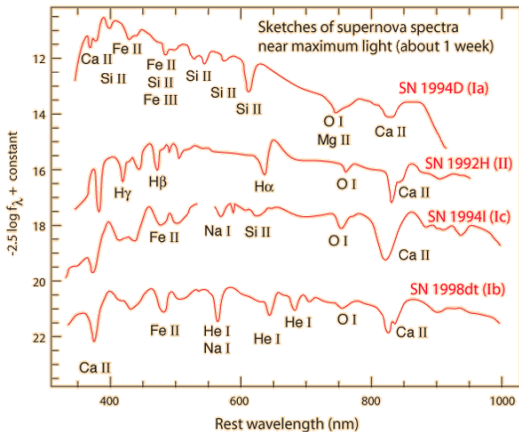
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Core collapse and
 $t \leq 3 \times 10^8$ years

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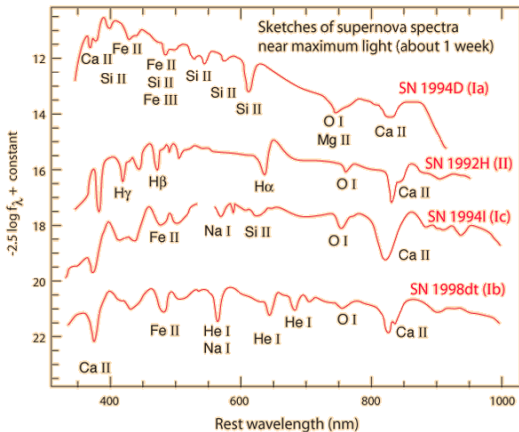
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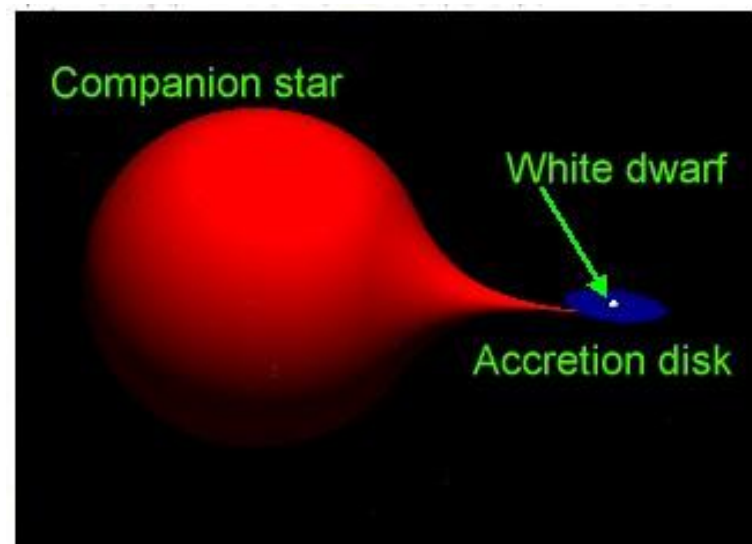


Core collapse and
 $t \leq 3 \times 10^8$ years

$t \sim 10^9$ years
 $M \leq 8 M_{\odot}$

Type Ia supernovae

White dwarf in a close binary system



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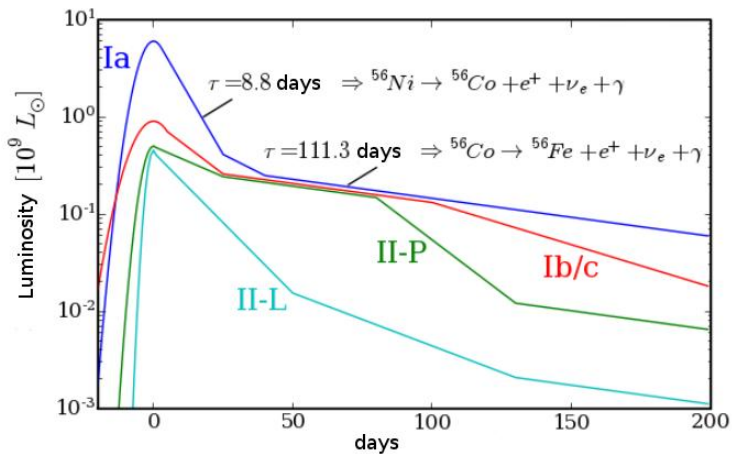
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Type Ia supernovae



SNLS-3yr: <http://www.physik.uni-bielefeld.de/igs/cosmology2009/guy.pdf> (Nov 29., 2009)

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- peak luminosity of SN Ia depends on the mass of ^{56}Ni ejected ($0.6 M_{\odot}$)

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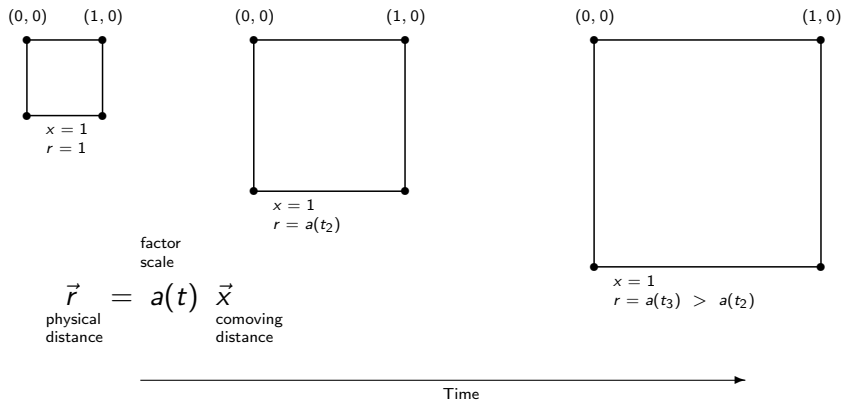
In 1990s two independent groups (SCP, HST) announced surprising discovery using SNe Ia as standard candles \rightarrow expansion of the Universe is accelerating (The results are consistent with some sort of 'dark energy', possibly Einstein's cosmological constant)

Data corrections

Uncorrected observations of SN Ia absolute magnitudes have an RMS scatter of $\sim 40\%$

- reduced via correlations between rate of decline, colour and intrinsic luminosity
- current standardization techniques normalize scatter to 15 % to 20 % (0.16 mag)
- Many popular methods (SALT(2), Stretch, (M)CLS(2k2), CMAGIC) enjoying similar level of success
- New method using flux ratios from a single flux-calibrated spectrum (0.12 mag) (arXiv:astro-ph/0905.0340)

Distances in cosmology



Friedmann equations

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 \text{Hubble rate} \rightarrow H^2(t) \equiv \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2} \leftarrow \text{Geometrical curvator} \\
 \frac{\partial \rho}{\partial t} + 3\frac{\dot{a}}{a}(\rho + p) = 0 \quad \begin{array}{l} k = 0 \rightarrow \text{flat universe} \\ k < 0 \rightarrow \text{open universe} \\ k > 0 \rightarrow \text{closed universe} \end{array}
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$\rho \leftarrow$ matter, dark energy, radiation, neutrinos.

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$$\Omega = \rho / \rho_{crit}$$

$$\left(\frac{H}{H_0}\right)^2 = \frac{\Omega_m}{a^3} + \frac{\Omega_k}{a^2} + \Omega_\Lambda$$

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Luminosity distance

The wavelength of light emitted from a receding object is stretched out so that the observed wavelength is larger than the emitted one. We define redshift (z) as

$$1 + z \equiv \frac{\lambda_{obs}}{\lambda_{emit}} = \frac{1}{a(t)}$$

It can be shown that the flux we observe is

$$F = \frac{L}{4\pi d_L^2}$$

Where we have introduced the luminosity distance d_L as

$$d_L = (1 + z) \int_0^z \frac{dz'}{H(z')} \quad (\text{in a flat universe})$$

Cosmological parameters from SNe Ia

Transforming luminosity to magnitude we get

$$m - M = 5 \log_{10}(d_L [\text{Mpc}]) + 25$$

We seek minimum of χ^2 statistics

$$\chi^2(H_0, \Omega_m, \Omega_\Lambda, w) = \sum_i \frac{(m_i - m(z_i; H_0, \Omega_m, \Omega_\Lambda, w))^2}{\sigma_{m_i}^2 + \sigma_m^2 + \sigma_v^2}$$

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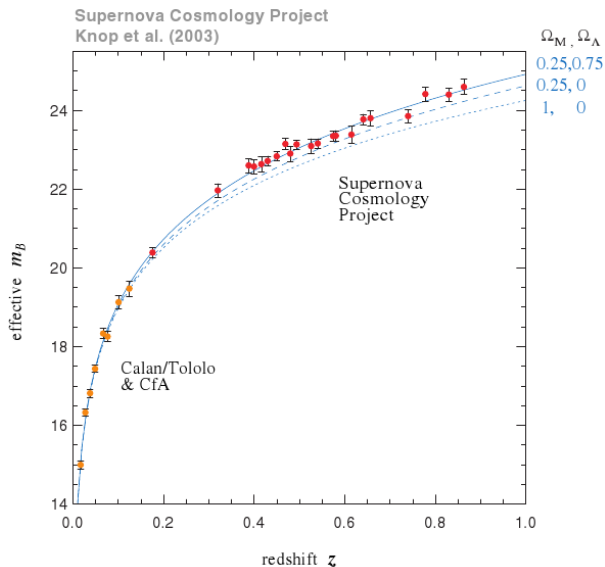
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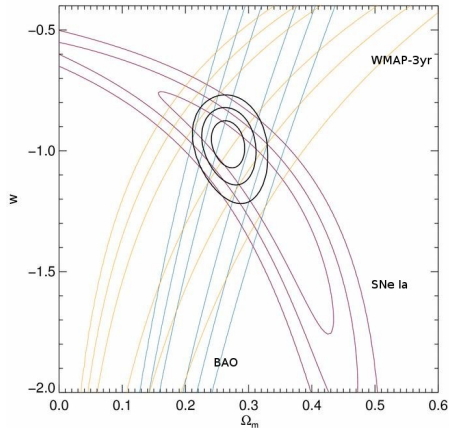
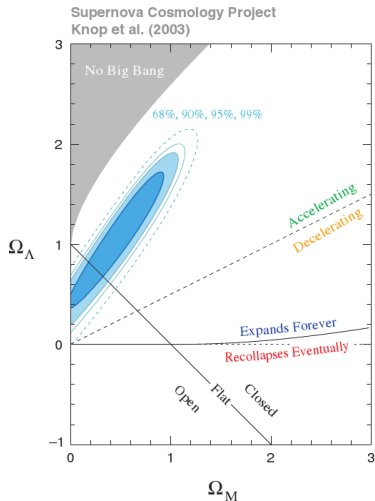
or maximum of probability L as

$$\log L \propto -\frac{1}{2}\chi^2 - \frac{1}{2} \sum_i \log(\sigma_{m_i}^2 + \sigma_m^2 + \sigma_v^2)$$

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http://128.111.9.106/online/snovae_c07/sullivan/pdf/Sullivan

<http://supernova.lbl.gov/> (Nov 29. 2009)

SnovaeConf KITP.pdf (Nov 29., 2009)



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$$\frac{\delta d_L}{d_L} = \frac{d_L^{me} - d_L^{th}}{d_L^{th}} = \hat{\mathbf{r}} \cdot \left(\mathbf{v} - \frac{(1+z)^2}{H(z)d_L} (\mathbf{v} - \mathbf{v}_o) \right)$$

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- the PV field can be derived ($\hat{\mathbf{r}} \cdot \mathbf{v}(\mathbf{r})$)

Correlation function #1

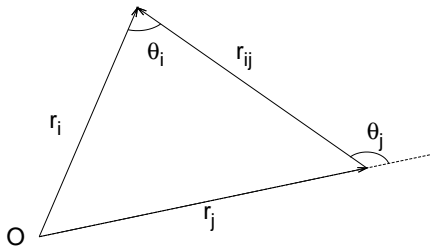
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- in linear theory in Fourier space $\mathbf{v}_{\mathbf{k}} \propto \frac{\delta(\mathbf{k}, z=0)}{k} \hat{\mathbf{k}}$
- the cross-correlation function for the matter density fluctuations is $\langle \delta(\mathbf{k}, z=0) \delta(\mathbf{k}', z=0) \rangle = 2\pi \delta^3(\mathbf{k} - \mathbf{k}') P(k, z=0)$

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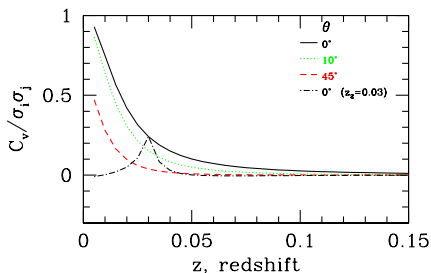
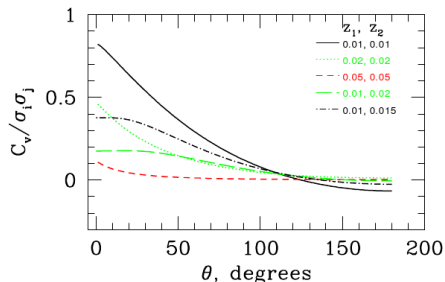
$$\xi(\mathbf{r}_i, \mathbf{r}_j) \propto \int_0^\infty \frac{dk}{2\pi^2} P(k) K(k|\mathbf{r}_{ij}|)$$

$$\langle (\mathbf{v}(\mathbf{r}) \cdot \hat{\mathbf{r}})^2 \rangle \propto \int_0^\infty P(k) dk$$

Correlation function #2

The PVs correlations lead to luminosity distance correlations and the covariance for the $\delta d_L/d_L$ for a pair of SN can be written as

$$C_v(i,j) = \left(1 - \frac{(1+z)^2}{H(z)d_L}\right)_i \left(1 - \frac{(1+z)^2}{H(z)d_L}\right)_j \xi(\mathbf{r}_i, \mathbf{r}_j)$$



Slosar et al, arXiv:astro-ph/0705.1718v5

Correlation function #3

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- alternative approach \rightarrow estimate the underlying density field from galaxy redshift surveys

Conclusion

- using SNe Ia data it has been proven that the expansion of the universe is accelerating
- indication of an unexpected gravitational physics → dark energy
- SNe Ia remain our best tool to constraint dark energy (precision and statistically significant samples)
- density inhomogeneities → deviation from Hubble flow: peculiar velocities
- in the limit of large data sets and small redshifts ($z \leq 0.1$) correlations between PVs contribute significantly to the overall error
- systematic uncertainties in relative distances due to: intervening dust, instrumental calibration and unknown evolution in z of SNe Ia luminosity ← limiting factors
- without PV correlations we measure parameters (H_0, Ω_m, w), with PV correlations we probe $P(k)$ and measure other parameters as well ($\Omega_b, n_s, \sigma_8, \dots$)